

Super-resolution, subspace methods, and Fourier matrices

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Outline

Mathematics of super-resolution

Non-harmonic Fourier matrices

Subspace methods

Fundamental limitations of super-resolution

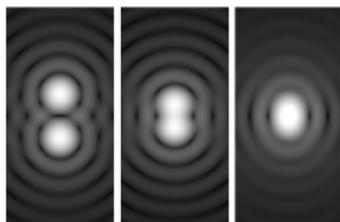
Multiple snapshot super-resolution

Conclusions

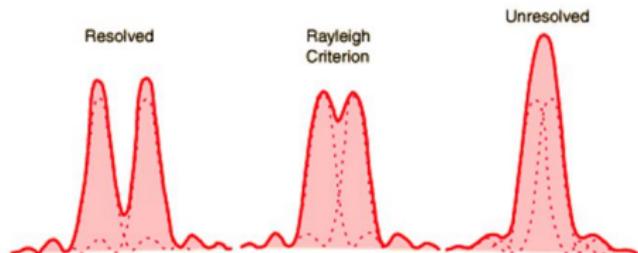
What is resolution?

The **Rayleigh length** of an imaging system is the smallest distance between two point sources that can be resolved.

Loosely referred to as the **resolution** of the imaging device.



(a) Airy disks



(b) One-dimensional plot

Resolution and Fourier transform

Spatial domain:

$$\left(\text{imaged object}\right) * \left(\text{point spread function}\right) = \left(\text{low-resolution image}\right)$$

Fourier domain:

$$\mathcal{F}\left(\text{imaged object}\right) \cdot \mathcal{F}\left(\text{point spread function}\right) = \mathcal{F}\left(\text{low-resolution data}\right)$$

Goal of super-resolution algorithms: Leverage prior information in order to overcome the inherent resolution limit of the imaging device – extract high-resolution features from observed low-resolution or coarse information.

Mathematical model of super-resolution [Donoho 1992]

Unknown: Atomic measure

$$\mu := \sum_{j=1}^S a_j \delta_{x_j}, \quad a_j \in \mathbb{C}, \quad x_j \in \mathbb{T} := [0, 1).$$

Known: Perturbation of M consecutive Fourier coefficients

$$y := \mathcal{F}_M \mu + \eta \in \mathbb{C}^M$$

where

$$\mathcal{F}_M \mu(m) := \{\widehat{\mu}(m)\}_{m=0}^{M-1}$$

and

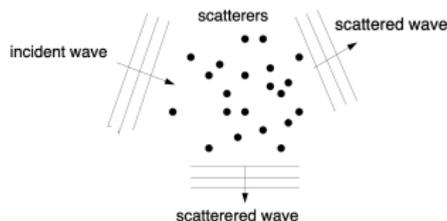
$$\widehat{\mu}(m) := \int_{\mathbb{T}} e^{-2\pi i m x} d\mu(x).$$

Goal: Recover μ . Primarily its support.

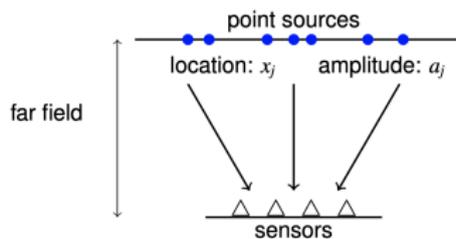
Imaging of point sources

Unknown: Atomic measure

$$\mu := \sum_{j=1}^S a_j \delta_{x_j}, \quad a_j \in \mathbb{C}, \quad x_j \in \mathbb{T} := [0, 1).$$



(a) Inverse scattering



(b) Direction of arrival

Applications: Super-resolution microscopy, geophysics, astronomy, remote sensing, inverse scattering, direction of arrival, and line spectral estimation.

Mathematical connections: Sparse recovery, non-harmonic Fourier analysis, sub-Nyquist sampling, and diffraction limited imaging.

Three main difficulties

1. It is a non-linear inverse problem:

$$\{(x_j, a_j)\}_{j=1}^S \leftrightarrow \mu \mapsto \mathcal{F}_M(\mu).$$

Well-posed in the sense that this map is injective if $M \geq 2S$ (Prony).

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2. We deal with non-harmonic sums:

$$\widehat{\mu}(m) = \sum_{j=1}^S a_j e^{-2\pi i m x_j}; \quad \widehat{\mu}(t) = \sum_{j=1}^S a_j e^{-2\pi i t x_j}.$$

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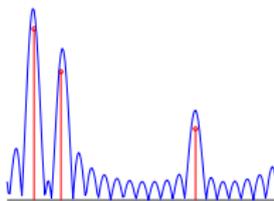
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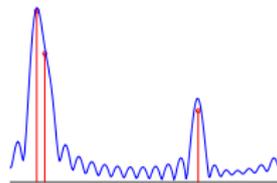
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3. Sources can be closely spaced:



(a) separation is $4/M$



(b) separation is $0.6/M$

Figure: $|\mu * D_M|$, where $D_M(x) = \frac{1}{M} \sum_{m=0}^{M-1} e^{2\pi i m x}$

Minimum separation:

$$\Delta := \min_{j \neq k} |x_j - x_k|_{\mathbb{T}}$$

Rayleigh length: $\approx 1/M$

Super-resolution factor: Rayleigh length divided by the target resolution

$$SRF := \frac{\text{standard resolution}}{\text{target resolution}} := \frac{1/M}{\Delta} = \frac{1}{M\Delta}.$$

Sub-Nyquist or diffraction limited regime:

$$SRF \geq 1 \quad \text{or equivalently} \quad \Delta \leq \frac{1}{M}.$$

Existing methods (up to around 2017–2018)

	$\Delta \gg 1/M$ or $SRF < 1$	$\Delta \ll 1/M$ or $SRF \gg 1$
1a. TV-min (μ complex)	$O(\ \eta\ _2)$	Can fail even if $\eta = 0$
1b. TV-min (μ positive)	$O(\ \eta\ _2)$	$O(\Delta^{-2S+1}\ \eta\ _2)$
2. Greedy methods	$O(\ \eta\ _2)$	Can fail if $\eta \neq 0$
3. Subspace methods	$O(\ \eta\ _2)$	Only numerical results
4. “Best” algorithm		$O(SRF^{2S-1}\ \eta\ _2)$

References:

1. TV-min: [Candès, Fernandez-Granda 2013, 2014], [Tang, Bhaskar, Shah, Recht 2013], [Duval, Peyré 2015], [Morgenshtern, Candès 2016], [Denoyelle, Duval, Peyré 2016], [Schiebinger, Robeva, Recht 2017]
2. Greedy: [Duarte, Baraniuk 2013], [Fannjiang, Liao 2012]
3. Subspace: [Liao, Fannjiang 2016], [Moitra 2015]
4. SR Limit: [Donoho 1992], [Demagnet, Nguyen 2015]

Main questions

Study the $\Delta \ll 1/M$ or $SRF \gg 1$ regime.

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Study the $\Delta \ll 1/M$ or $SRF \gg 1$ regime.

1. **Computational:** Are there accurate and efficient algorithms?
2. **Computational analysis:** What are the performance guarantees of said algorithms?
3. **Information theoretic:** What is the best possible recovery rate independent of algorithm?

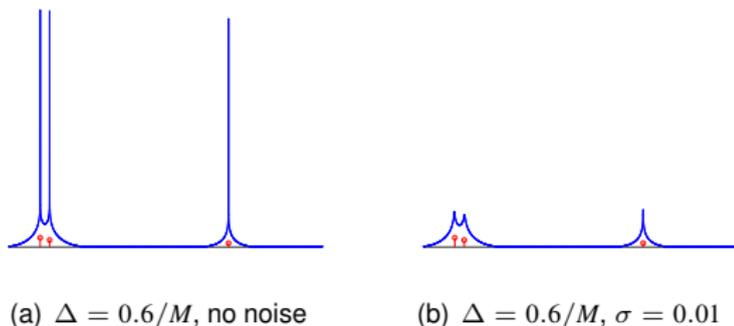
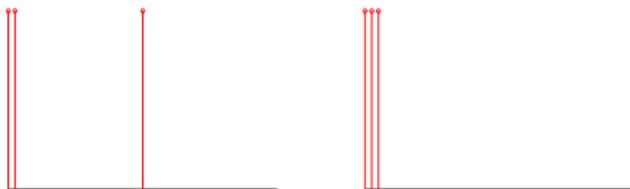
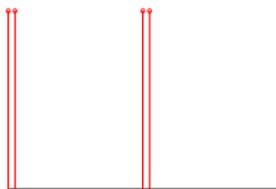
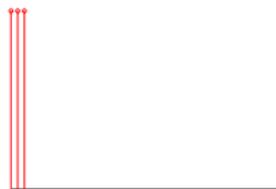
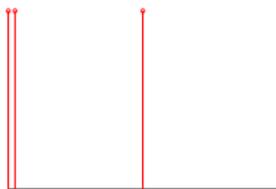


Figure: Output of the MUSIC algorithm

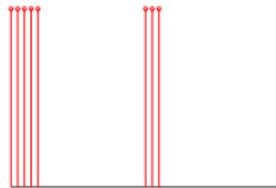
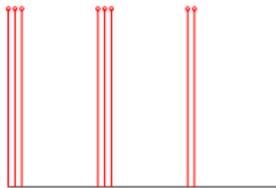
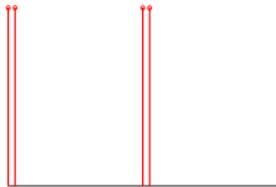
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Fourier matrices

Unknown: Atomic measure $\mu := \sum_{j=1}^S a_j \delta_{x_j}$, where $a_j \in \mathbb{C}$ and $x_j \in \mathbb{T} := [0, 1)$.

Known: Consecutive M noisy Fourier coefficients

$$y := \mathcal{F}_M \mu + \eta, \quad \mathcal{F}_M \mu(m) := \int_{\mathbb{T}} e^{-2\pi i m x} d\mu(x) \quad \text{for } m = 0, 1, \dots, M-1.$$

Fourier matrix associated with $X = \{x_j\}_{j=1}^S$:

$$\Phi_M := \Phi_M(X) := \begin{bmatrix} 1 & 1 & \dots & 1 \\ e^{-2\pi i x_1} & e^{-2\pi i x_2} & \dots & e^{-2\pi i x_S} \\ \vdots & \vdots & & \vdots \\ e^{-2\pi i (M-1)x_1} & e^{-2\pi i (M-1)x_2} & \dots & e^{-2\pi i (M-1)x_S} \end{bmatrix},$$

and so

$$y = \Phi_M(X)a + \eta.$$

Fourier matrices and super-resolution

Fourier matrix associated with $X = \{x_j\}_{j=1}^S$:

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Minimum singular value $\sigma_S(\Phi_M(X))$ controls:

- ▶ Robustness of subspace methods (will return to this later).
- ▶ Min-max error or information theoretic limit (will return to later).

Connections between super-resolution and Fourier matrices is hinted at by earlier work [Donoho 1992], [Demanet, Nguyen 2015], [Moitra 2015].

Known results about Fourier matrices

Dichotomy:

- ▶ If $\Delta \geq C/M$ for some $C > 1$, then $\sigma_S(\Phi_M) \gtrsim \sqrt{M}$.

[Vaaler 1985], [Moitra 2015], based on work by Beurling and Selberg.

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For example, [Gautschi 1962] showed that for a square general Vandermonde matrix with nodes $Z = \{z_1, \dots, z_S\} \subseteq \mathbb{C}$,

$$\sigma_S(\Phi_S(Z)) \geq \min_{1 \leq j \leq S} \prod_{k=1, k \neq j}^S \frac{|z_k - z_j|}{1 + |z_k|},$$

where equality holds if and only if z_1, \dots, z_S lie on the same line in \mathbb{C} .

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This bound is sharp, but can be improved under certain configurations.

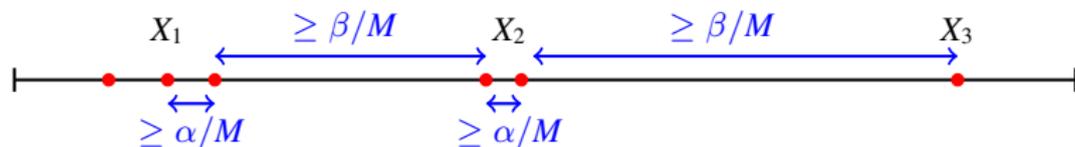
Separated clumps model

Separated clumps: (Two-scale model) A set $X = \{x_j\}_{j=1}^S$ consists of separated clumps with parameters (R, M, S, α, β) if we have the disjoint union

$$X = \bigcup_{r=1}^R X_r$$

where each X_r is contained in an interval of length $1/M$ and

1. (Intra-clump separation) $\Delta \geq \alpha/M$ where $\alpha \leq 1$.
2. (Inter-clump separation) If $R > 1$, then $\text{dist}(X_j, X_k) \geq \beta/M$ where $\beta \geq 1$.



Two Extremes:

- ▶ $R = S$: each clump is a single point and $\Delta \geq 1/M$.
- ▶ $R = 1$: one clump contains all S points.

Lower bound for separated clumps

Theorem (L., W. Liao, ACHA 2021)

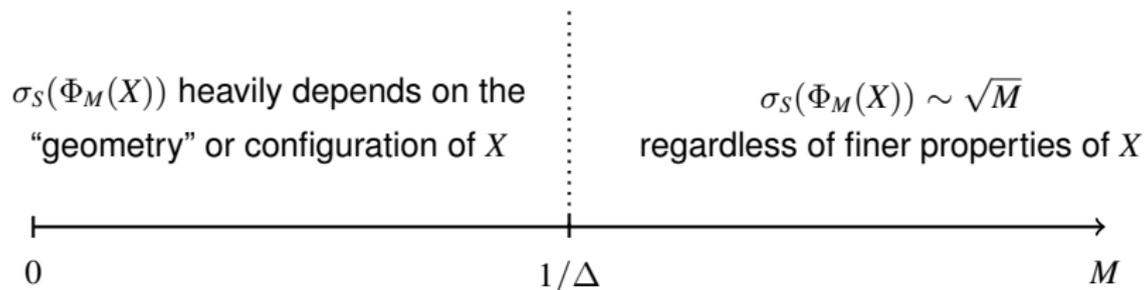
Suppose X consists of separated clumps with parameters (R, M, S, α, β) , where $M \geq S^2$, $\alpha \in (0, 1)$, and $\beta \geq \max_r 20\sqrt{S} \lambda_r^{5/2} / \sqrt{\alpha}$ where $\lambda_r = |X_r|$. Then there exist explicit constants $\{C_r\}_{r=1}^R$ that do not depend on M and α such that

$$\sigma_S(\Phi_M(X)) \geq \sqrt{M} \left(\sum_{r=1}^R C_r^2 \left(\frac{1}{\alpha} \right)^{2\lambda_r - 2} \right)^{-1/2}.$$

Rough estimate: Letting $\lambda = \max_r \lambda_r$, note that $SRF = \frac{1/M}{\alpha/M} = 1/\alpha$, so

$$\sigma_S(\Phi_M(X)) \geq C(\lambda) \sqrt{\frac{M}{R}} SRF^{-\lambda+1} = C(\lambda) \sqrt{\frac{M}{R}} (M\Delta)^{\lambda-1}.$$

Phase transition:



Theorem (L., W. Liao, ACHA 2021)

If X consists of separated clumps with parameters (R, M, S, α, β) where $M \geq S^2$, $\alpha \in (0, 1)$, and $\beta \geq \max_r 20\sqrt{S} \lambda_r^{5/2} / \sqrt{\alpha}$ where $\lambda_r = |X_r|$, there exist explicit constants $\{C_r\}_{r=1}^R$ where C_r only depends on λ_r and

$$\sigma_S(\Phi_M(X)) \geq \sqrt{M} \left(\sum_{r=1}^R C_r^2 \left(\frac{1}{\alpha} \right)^{2\lambda_r - 2} \right)^{-1/2}.$$

Brief history on “geometrically” aware lower bounds on $\sigma_S(\Phi_M)$:

- ▶ For general Vandermonde matrices, see [Gautschi 1962] for square case, and [Bazán 2000] for rectangular ones.
- ▶ For rectangular Fourier matrices, this result appeared in 2018, concurrently with [Batenkov, Demanet, Goldman, Yomdin, 2020].
- ▶ Follow up papers [Kunis, Nagel, 2020], [Demanet, Goldman, Yomdin, 2021].

Proof technique

Duality: If $v \in \mathbb{C}^S$ is the S -th right singular vector of $\Phi_M(X)$, then

$$\sigma_S(\Phi_M(X)) = \max_{\substack{f \text{ trig poly degree } M-1 \\ f(X)=v}} \|f\|_{L^2(\mathbb{T})}^{-1}.$$

Method: Construct trigonometric polynomials $\{L_{X,j}\}_{j=1}^S$

1. Each $L_{X,j}$ has degree $M - 1$
2. $L_{X,j}(x_j) = 1$
3. $L_{X,j}(x_k) = 0$ for all x_k in the same clump as x_j
4. $L_{X,j}$ decays quickly away from x_j .

$$\left\| \underbrace{\sum_{j=1}^S v_j L_{X,j}}_{\text{approximate interpolant}} \right\|_{L^2(\mathbb{T})} \leq \left(\sum_{j=1}^S \|L_{X,j}\|_{L^2(\mathbb{T})}^2 \right)^{1/2}.$$

By a robust version of duality,

$$\sigma_S(\Phi_M(X)) \gtrsim \left(\sum_{j=1}^S \|L_{X,j}\|_{L^2(\mathbb{T})}^2 \right)^{-1/2}$$

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Subspace (algebraic) methods:

1. MUSIC [Schmidt 1986]
2. ESPRIT [Roy, Kailath 1989]
3. MPM [Hua, Sarkar 1990]

Hankel matrix: Suppose $M \geq 2S + 1$ and $L \approx M/2$.

$$H(y) := \begin{bmatrix} y_0 & y_1 & \cdots & y_{M-1-L} \\ y_1 & y_2 & \cdots & y_{M-L} \\ \vdots & \vdots & & \vdots \\ y_L & y_{L+1} & \cdots & y_{M-1} \end{bmatrix}.$$

Fourier decomposition of Hankel matrix

$$H(\mathcal{F}_M \mu) = \Phi_L(X) \text{diag}(a) \Phi_{M-1-L}(X)^t.$$

Which implies

$$H(y) = H(\mathcal{F}_M \mu) + H(\eta).$$

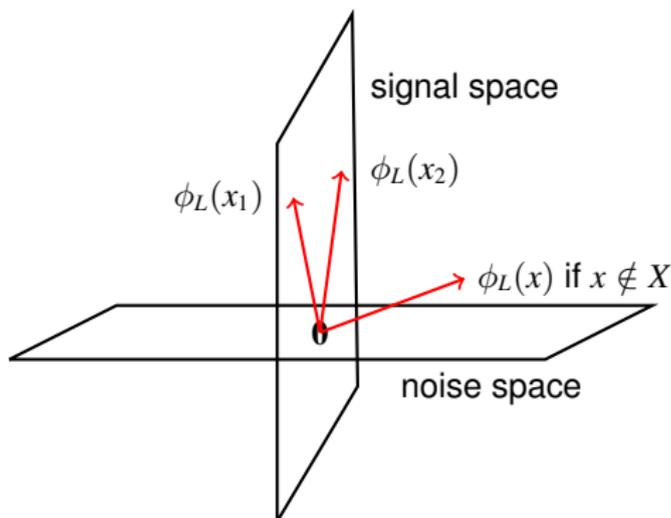
Subspace methods continued

Compute the best rank S approximation of $H(y)$ by SVD truncation:

$$\tilde{U}\tilde{\Sigma}\tilde{V}^* = \text{best rank } S \text{ approximation of } H(y).$$

If the noise is sufficiently small,

$$\text{Range}(\tilde{U}) \approx \text{Range}(\Phi_L(X))$$



Estimation of signal parameters via rotational invariance techniques

ESPRIT algorithm:

▶ **Noiseless case:**

1. Compute a matrix U whose columns form an orthonormal basis for the range of $\Phi_M(X)$.
2. Let U_0 and U_1 be the first and last L rows of U , respectively.
3. The eigenvalues of $U_0^\dagger U_1$ are $\{e^{2\pi i x_j}\}_{j=1}^S$, from which we extract $\{x_j\}_{j=1}^S$.

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Subspace methods require knowing S , or one estimates S through some other knowledge such as noise level. Recent work [P. Liu, H. Zhang 2021]

Theorem (L., W. Liao, A. Fannjiang, IEEE IT 2020)

For any S -atomic μ with support X and any η such that $\|H(\eta)\|_2$ is sufficiently small, if \tilde{X} is the output of ESPRIT, we have

$$md(X, \tilde{X}) \leq \frac{C(M, S)}{\underbrace{(\min_j |a_j|) \sigma_S^2(\Phi_L(X))}_{\text{numerical conditioning}}} \|\eta\|_2.$$

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Under the clumps model (under additional assumptions on X),

$$md(X, \tilde{X}) \leq \underbrace{\frac{C(R, M, S, \lambda) SRF^{2\lambda-2}}{\min_j |a_j|}}_{\text{numerical conditioning}} \|\eta\|_2.$$

Comparison: Previous best bounds for ESPRIT are

$$O(\|\eta\|_2 \sigma_S^{-5}(\Phi_{M/2})) \quad [\text{Aubel, Bölcskei, 2016}],$$

$$O(\|H(\eta)\|_2 \sigma_S^{-4}(\Phi_{M/2})) \quad [\text{Fannjiang, 2016}].$$

ESPRIT is near optimal: Upper bound is sharp in terms of SRF , and is one M factor away from min-max lower bound in [Batenkov, Goldman, Yomdin, 2020].

ESPRIT automatically adapts to the geometry: The algorithm does not require knowledge of the clump parameters!

Non-harmonic uncertainty principle

Concentration: Given a S -atomic measure μ and $L \geq S$, the quantity

$$C_L(\mu) := \frac{|\widehat{\mu}(0)|^2}{\sum_{m=0}^L |\widehat{\mu}(m)|^2}.$$

Uncertainty principle: If μ is S -atomic, expect its Fourier coefficients to not be perfectly localized.

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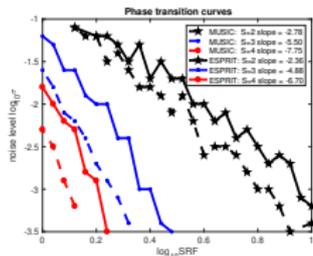
If $L \geq S > 1$, then

$$C_{L,S} := \sup_{\substack{\mu \text{ is } S\text{-atomic} \\ \mu \neq 0 \\ \mu \text{ complex}}} C_L(\mu) \leq 1 - 4^{-S},$$

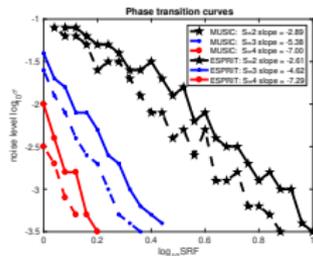
$$C_{L,S,\mathbb{R}} := \sup_{\substack{\mu \text{ is } S\text{-atomic} \\ \mu \neq 0 \\ \mu \text{ real}}} C_L(\mu) \leq 1 - (8S - 1)^{-1}.$$

Can be seen as a quantitative version of a result in [Donoho, Stark 1989].

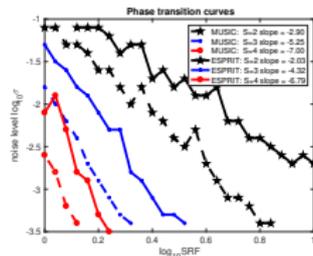
Numerical simulations



(a) 1 clump



(b) 2 clumps



(c) 3 clumps

Figure: The phase transition curves below which the algorithm succeeds (defined to be successful if $\text{md}(X, \tilde{X}) < \Delta/2$) with probability at least 95% for $\lambda = 2, 3, 4$ with respect to $\log_{10}(SRF)$ (x-axis) and $\log_{10} \sigma$ (y-axis). The slopes are computed by least squares.

Numerical simulations continued

	$\lambda = 2$	$\lambda = 3$	$\lambda = 4$	Numerical	Theoretical
1-clump: MUSIC	2.78	5.50	7.75	$2.49\lambda - 2.11$	$2\lambda - 2$
2-clump: MUSIC	2.89	5.38	7.00	$2.06\lambda - 1.08$	$2\lambda - 2$
3-clump: MUSIC	2.90	5.25	7.00	$2.05\lambda - 1.10$	$2\lambda - 2$
4-clump: MUSIC	3.01	5.12	8.50	$2.75\lambda - 2.70$	$2\lambda - 2$
1-clump: ESPRIT	2.36	4.88	6.70	$2.17\lambda - 1.86$	$2\lambda - 2$
2-clump: ESPRIT	2.61	4.62	7.29	$2.34\lambda - 2.18$	$2\lambda - 2$
3-clump: ESPRIT	2.03	4.32	6.79	$2.38\lambda - 2.76$	$2\lambda - 2$
4-clump: ESPRIT	1.81	4.34	6.43	$2.31\lambda - 2.74$	$2\lambda - 2$

Table: Slopes extracted from the previous phase transition curves of MUSIC and ESPRIT.

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Min-max error for sparse model on the grid

Grid model: Suppose $X \subseteq \{\frac{n}{N}\}_{n=0}^{N-1}$ for some large N . Every S -atomic μ can be identified with a S -sparse vector $u \in \mathbb{C}_S^N$.

Min-max error: Accuracy of the “best” possible algorithm (including those that have exponential run-time),

$$E(M, N, S, \delta) = \inf_{\substack{\phi: y \rightarrow \mathbb{C}_S^N \\ y = \mathcal{F}_M u + \eta}} \sup_{\substack{u \in \mathbb{C}_S^N: \|u\|_2 \leq 1 \\ \eta \in \mathbb{C}^M: \|\eta\|_2 \leq \delta}} \|\phi(u, \eta) - u\|_2.$$

No algorithm can beat the min-max error, and in particular,

ESPRIT error for this sparsity model $\geq E(M, N, S, \delta)$.

Sharp estimate on the min-max error

Theorem (L., W. Liao, ACHA 2021)

Let $S \geq 1$ and $M \geq 4S$. For N sufficiently large,

$$\frac{1}{\sqrt{M}} SRF^{2S-1} \delta \lesssim_{M,S} E(M, N, S, \delta) \lesssim_{M,S} \frac{1}{\sqrt{M}} SRF^{2S-1} \delta,$$

where $SRF = N/M$.

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where $SRF = N/M$.

Remarks:

1. Non-asymptotic bounds for the min-max error with explicit constants.
2. Single clump with S points is the worst case scenario?
3. Proof relies on showing that $E(M, N, S, \delta)$ is related to $\sigma_S(\Phi_M(X \cup \tilde{X}))$, where $X \cup \tilde{X}$ is a set of cardinality at most $2S$.

Related work: [Donoho 1992], [Demagnet, Nguyen 2015], [Batenkov, Goldman, Yomdin, 2020]

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How to incorporate additional information?

Unknown: Time-dependent atomic measure, where the locations are fixed, but the amplitudes are time varying

$$\mu_t := \sum_{j=1}^S a_j(t) \delta_{x_j}, \quad a_j(t) \in \mathbb{C}, \quad x_j \in \mathbb{T} := [0, 1).$$

Known: Perturbed consecutive M Fourier coefficients at times t_1, \dots, t_n ,

$$y(t_\ell) := \mathcal{F}_M(\mu_{t_\ell}) + \eta(t_\ell), \quad \ell = 1, \dots, n.$$

This is the **multiple-snapshot** problem, in contrast to the **single-snapshot** version earlier.

Multiple snapshot ESPRIT

ESPRIT readily extends to the multi-snapshot case:

1. Empirical covariance matrix

$$Y := \frac{1}{n} \sum_{\ell=1}^n y(t_\ell) y(t_\ell)^*.$$

2. \tilde{U} is the best rank S approximation of Y .
3. Find the eigenvalues of $\tilde{U}_0^\dagger \tilde{U}_1$, project to the unit complex circle, and extract their arguments.

What is the performance of ESPRIT

Is there any advantage of taking Fourier measurements over time?

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multi-snapshot error = single-snapshot error.

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This comes from just doing naive averaging:

$$\text{avg}(y(t_1), \dots, y(t_n)) = \mathcal{F}_M \left(\sum_{j=1}^S \text{avg}(a_j(t_1), \dots, a_j(t_n)) \delta_{x_j} \right) + \underbrace{\frac{1}{n} \sum_{\ell=1}^n \eta(t_\ell)}_{\sqrt{n} \text{ cancellation}}$$

Theorem (L., Z. Zhu, W. Gao, W. Liao, preprint 2021)

1. Assume that each $\eta(t_\ell) \sim CN(0, \nu^2 I)$ (this can be relaxed significantly).
2. Assume that $a(t_1), \dots, a(t_n)$ span \mathbb{C}^S . Or equivalently the amplitude covariance matrix

$$A = \frac{1}{n} \sum_{\ell=1}^n a(t_\ell) a(t_\ell)^*$$

is strictly positive definite.

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is strictly positive definite.

For n sufficiently large, the output \tilde{X} of ESPRIT satisfies, with probability at least $1 - \delta$,

$$md(X, \tilde{X}) \leq \underbrace{\frac{C(M, S)}{\sqrt{n} \sqrt{\lambda_S(A)} \sigma_S(\Phi_M)}}_{\text{numerical conditioning}} \nu \left(1 + \sqrt{\frac{\log(1/\delta)}{M}} \right).$$

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Under the clumps model,

$$md(X, \tilde{X}) \leq \underbrace{\frac{C(M, S) \text{SRF}^{\lambda-1}}{\sqrt{n} \sqrt{\lambda_S(A)}}}_{\text{numerical conditioning}} \nu \left(1 + \sqrt{\frac{\log(1/\delta)}{M}} \right).$$

Single and multi snapshot

► **Single snapshot:**

$$\text{ESPRIT error} \lesssim_{M,S,a} \frac{\text{noise}}{\sigma_S^2(\Phi_M)} \lesssim_{M,S,a} SRF^{2\lambda-2} \text{noise}.$$

Dependence on SRF and noise match the min-max rate, so ESPRIT is optimal in this sense.

► **Multi snapshot:**

$$\text{ESPRIT error} \lesssim_{M,S,a} \frac{\text{noise}}{\sqrt{n}\sigma_S(\Phi_M)} \lesssim_{M,S,a} \frac{1}{\sqrt{n}} SRF^{\lambda-1} \text{noise}.$$

Dependence on SRF , number of snapshots n , and noise match a Cramer-Rao lower bound [L., Zhu, Gao, Liao 2021], so ESPRIT is optimal in this sense.

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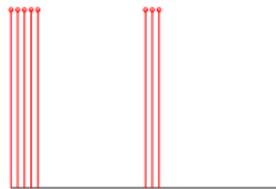
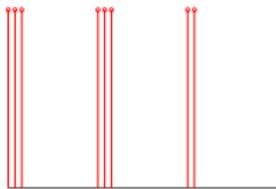
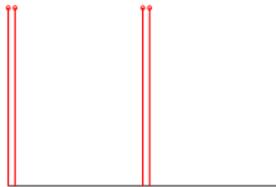
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Difficulty of super-resolution

Order from least to most challenging in the limit $\Delta \rightarrow 0$:



(a) $O(SRF^2 \|\eta\|_2)$



(b) $O(SRF^4 \|\eta\|_2)$



(c) $O(SRF^8 \|\eta\|_2)$

Some remarks

- ▶ Classical work by Rayleigh and Abbe study the simplest scenario where there are only two point sources. Our results give a description for more complicated arrangements.

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- ▶ Our results show that super-resolution is possible in more situations than previously thought, since the error is $O(SRF^{2\lambda-2}\|\eta\|_2)$ as opposed to $O(SRF^{2S-2}\|\eta\|_2)$.

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- ▶ ESPRIT is provably optimal (in terms of SRF), so criticism about “instability of subspace methods” is not fair.
- ▶ Since $\sigma_S(\Phi_M(X)) > 0$ for any set X , an error bound of the form

$$\text{error} \lesssim \frac{\text{noise}}{\sigma_S^2(\Phi_M(X))}$$

means that (at least for point sources and Fourier measurements) separation is never the real bottleneck of super-resolution and that noise is the real culprit.

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Thank you!